Cosmic shear full nulling: sorting out dynamics, geometry and systematics

MNRAS 2014, FB, Nishimichi, Taruya

CEA IAP meeting 13 novembre 2015

Growth of structure in a self-gravitating dust fluid



Reconstruction of the power spectrum: from sPT to Multi-point propagator reconstruction



1 [1 14 -1]

Good performance of PT calculations (up to 2 loops)

in real space



Power spectra up to 1-loop and 2-loop order

Taruya, FB, Nishimichi, Codis '12 Crocce, Scocimarro, FB, '12 Ist computation of 2-loop order effects in Okamura, Taruya, Matsubara, '11



 Public codes for fast computations of power spectra at 2-loop order are now available.

> http:// maia.ice.cat/ crocce/ mptbreeze/ http://wwwutap.phys.s.utokyo.ac.jp/ ~ataruya/ regpt_code.html

•Theoretical predictions are within 1% accuracy.

Charting PT in Fourier space

 Accuracy of linear and 1-loop results is obtained with a comparison with 2-loop results;

 Accuracy of 2-loop results is obtained with a comparisons with different prescriptions (that depart at 3-loop order)



The nulling

 the idea is to exploit nulling in order to avoid linear-nonlinear mixing in the projection effects.

MNRAS 2014, FB, Nishimichi, Taruya

Mixture of information

$$P_{\kappa}(\ell) = \frac{9H_0^4 \Omega_{\mathrm{m}}^2}{4c^4} \int_0^{\chi_{\mathrm{h}}} \mathrm{d}\chi \frac{g^2(\chi)}{a^2(\chi)} P_{\delta}\Big(\frac{\ell}{f_{\mathrm{K}}(\chi)}; z(\chi)\Big),$$

- Nonlinearity is more important on larger k at smaller z.
- Linear information at high z is contaminated by nonlinearity at low z.
- For any *l*, nearby nonlinear structure always comes in to the angular power spectrum.



Tomography?

$$P_{\kappa}^{(ij)}(\ell) = \frac{9H_0^4 \Omega_{\rm m}^2}{4c^4} \int_0^{\chi_{\rm h}} \mathrm{d}\chi \frac{g^{(i)}(\chi)g^{(j)}(\chi)}{a^2(\chi)} P_{\delta}\Big(\frac{\ell}{f_{\rm K}(\chi)}; z(\chi)\Big)$$

More information is accessible by tomography.

- can probe the structure growth at different times
- Lensing profiles overlap with each other.
- All are strongly affected by nearby structure.



Cosmic shear full nulling

- Can we control the lensing response function as desired?
 - non-zero response only for a given interval of χ (or z)?
 - Let us introduce a weight function $w(\chi_s)$, which we can set arbitrarily.

condition



Cosmic shear full nulling (cont.)

 The condition for the weighted source number density, wns, can be rewritten as

$$\int \mathrm{d}\chi_{\mathbf{s}} w(\chi_{\mathbf{s}}) n_{\mathbf{s}}(\chi_{\mathbf{s}}) = 0,$$

$$\int d\chi_{s} \frac{w(\chi_{s})n_{s}(\chi_{s})}{\tan(\sqrt{K}\chi_{s})} = 0, \qquad (K > 0),$$
$$\int d\chi_{s} \frac{w(\chi_{s})n_{s}(\chi_{s})}{\chi_{s}} = 0, \qquad (K = 0),$$
$$\int d\chi_{s} \frac{w(\chi_{s})n_{s}(\chi_{s})}{\tanh(\sqrt{-K}\chi_{s})} = 0, \qquad (K < 0).$$

We can always find solutions when the sources have a distribution with a finite width!

Demonstration: discrete source distribution

 One can always find a solution with 3 distinct source redshifts
 @ z = z₁, z₂ & z₃ (z₁ < z₂ < z₃).

$$\begin{pmatrix} wn_{s})_{1} = \chi_{1}(\chi_{3} - \chi_{2}), \\ (wn_{s})_{2} = -\chi_{2}(\chi_{3} - \chi_{1}), \\ (wn_{s})_{3} = \chi_{3}(\chi_{2} - \chi_{1}), \end{cases}$$
 (K=0)
$$\begin{pmatrix} \tilde{\kappa}_{1} \\ \tilde{\kappa}_{2} \\ \tilde{\kappa}_{3} \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ -1 & 1 & 0 \\ (wn_{s})_{1} & (wn_{s})_{2} & (wn_{s})_{3} \end{pmatrix} \begin{pmatrix} \kappa_{1} \\ \kappa_{2} \\ \kappa_{3} \end{pmatrix}$$



Demonstration: discrete source distribution

We can play the same game with more source planes...

$$\tilde{\kappa}_{i} = \sum_{j} \mathcal{W}_{ij} \kappa_{j},$$

$$\mathcal{W}_{ij} = \begin{bmatrix} 1 & 0 & 0 & 0 & 0 & 0 \\ -1 & 1 & 0 & 0 & 0 & 0 \\ n_{132} & n_{213} & n_{321} & 0 & 0 & 0 \\ 0 & n_{243} & n_{324} & n_{432} & 0 & 0 \\ 0 & 0 & n_{354} & n_{435} & n_{543} & 0 \\ 0 & 0 & 0 & n_{465} & n_{546} & n_{654} \end{bmatrix}$$



Nulling: it is possible to choose the

k contributions for I: 200, 300, 400, 500





FIG. 4: Adopted lens distribution for the available redshift source planes







FIG. 4: Adopted lens distribution for the available redshift source planes



A nulling approach, continuous limit



Figure 7. The adopted profiles in redshifts for the sources (top panel) and the resulting profiles for the lens distribution (bottom panel) for the fiducial cosmological model. The profiles for the sources have been obtained from the form (26) after they are convolved with a kernel that mimic the photometric redshifts error distribution.

Charting PT

- To get insights into the development of gravitational instabilities;
- To test/complement N-body simulations;
- Provide predictions from first principles in a large variety of models, and for a large numbers of parameters.

Charting PT in harmonic space (for cosmic shear observations)



ł

Conclusions

There has been a lot of developments beyond mere recipes for the constructions of power spectra beyond linear regime.

The goal is to obtain **controllable** predictions from **first principles**.

These results are directly applicable to cosmic shear observations taking advantage of the nulling construction.